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A Process for Synthesizing Bandlimited Chaotic Waveforms for Digital Signal Transmission

Chance M. Glenn

Abstract--In our development of a chaotic oscillator technology to produce high-quality communication signals, we have found a novel method for limiting the out-of-band spectral power from chaotic oscillators. This development is an important breakthrough that has allowed us to make a major step toward a commercially viable technology.

Index Terms--chaos, chaotic, communications, transmission, digital, wireless, bandlimiting, bandwidth, filtering, symbolic dynamics, binary

I. INTRODUCTION

In a practical wireless communication system, the signal transmitted from the antenna must be limited to a finite range of frequencies. Such signals are known as bandlimited signals, and if the range of frequencies does not include dc (zero frequency) they are called passband signals. Typically, passband signals are bandlimited signals that occupy a small percentage bandwidth about a center or carrier frequency. Multiple signals can then be transmitted by using passbands that do not overlap, then separating the signals at the receivers by filtering all but the passband of interest. This method of sending multiple signals, known as frequency-division multiplexing, is the basis for many multiple user systems in use today. When many users can access a system on either a fixed or flexible frequency division plan, the method is known as frequency-division multiple access. Although some methods, such as CDMA (code division multiple access) do not rely upon frequency division for signal separation, the signals are still limited to a passband defined by the FCC.1

If the signals generated by chaotic systems are to be used in commercially viable systems, then the transmitted signal must be bandlimited in some way. One obvious and often-used way to do this is to generate a signal that has considerable out-of-band spectral power, and then filter the signal to remove most of the power that lies outside of the desired frequency band. This method is, however, costly in many ways. Filters that can remove out-of-band energy well while passing the signal in band with minimal distortion are expensive and take up space. The power outside of the band is also wasted, and must be dissipated to heat. This wasted power translates to a higher required transmitter power and faster battery drain.

In our development effort to produce commercially viable chaotic oscillator technology, we have discovered a particularly simple and effective way to limit the out-of-band radiation from a chaotic oscillator. This method causes the oscillator itself to produce a signal with bandwidth-constrained signal power, and thus there is no need for a filter or waste of power. In doing so we are applying a principle that may be used for more general signal shaping or spectral shaping.

The idea behind our bandlimited chaotic oscillation (BCO) synthesis method is based on our segment hopping method of oscillator control. In segment hopping, a digital source produces an analog waveform that is used to guide the transmit oscillator. The guide signal is an analog copy of a signal that could be produced by the transmit oscillator itself, except that it follows a pre-defined symbol sequence that contains the digital information being transmitted. In this scheme, the transmit oscillator is acting as an amplifier for the guide signal, because the output of the transmitter can be much higher in power than the power drawn from the guide source.

II. SYNTHESIS/ANALYSIS

A prototypical chaotic oscillator used to produce signals useful for digital communication is the Lorenz system. It was the oscillator first used by Hayes to introduce the notion of controlling symbolic dynamics, a process to encode digital information in the oscillations of a chaotic system. The Lorenz system is described by a three dimensional system of equations having the form:

\[ \dot{x} = \sigma y - \sigma x \]
\[ \dot{y} = \rho x - y - xz \]
\[ \dot{z} = xy - \beta z \]

where \( \sigma, \rho \) and \( \beta \) are parameters that Lorenz originally set to 10, 28, and 8/3 respectively. The state-space attractor defined by these equations takes on a double-lobed structure which lends itself nicely to a binary symbol partition. Fig. 1 is an example of a two-dimensional projection of the solutions of the Lorenz equations, showing...
state-variables \( x \) and \( y \). Two Poincaré surfaces are placed in positions piercing the focal points and extending outward beyond the boundaries of attractor.

\[ \begin{align*}
\text{Figure 1.} \quad & \text{Two-dimensional projection of the solutions to the Lorenz equations showing state coordinates } x \text{ and } y \text{ and the binary partitioning of the state-space.} \\
\text{Figure 2.} \quad & \text{Time dependent solution of the Lorenz equations for state coordinate } x(t). \\
\end{align*} \]

The time varying state-variable, \( x(t) \), is a bipolar waveform well-suited for baseband transmission of digital information. Fig. 2 is an example of the waveform in it’s generalized time coordinates, and Fig. 3 a plot of the frequency content of this signal with respect to it’s average cycle frequency. The average cycle frequency is the reciprocal of the average cycle time. The cycle time is defined as the time it takes a point on the attractor to travel from one Poincaré surface to the other, or back to itself. The average cycle time is the mean time, in dimensionless units, calculated by integrating the system with a fixed integration step and collecting thousands of surface crossings. \( T_{avg} \) was found to be about 1.7, with a variance of about 1.8.

\[ \begin{align*}
\text{Figure 3.} \quad & \text{Frequency content of the state coordinate } x(t) \text{ with respect to the average cycle frequency.} \\
\end{align*} \]

Fig. 4 is a block diagram of the implementation of a bandlimited chaotic signal source. A binary sequence is fed into a discrete-time, bandlimited, segment-hopping source. We will describe this source in the next section. The output signal, \( y[n;t] \), is a sampled waveform used as a guide signal to synchronize a continuous-time Lorenz oscillator to it. A single state-variable synchronization method is used to lock the oscillator to the dynamics, thus to the embedded digital sequence of the guide signal.

\[ \begin{align*}
\text{Figure 4.} \quad & \text{Block diagram of the implementation of a bandlimited digital signal synthesizer.} \\
\end{align*} \]

Since Pecora and Carroll first introduced the concept of synchronization of chaotic oscillator circuits\(^5\) many methods have been developed and realized. The following mathematical model works extremely well:

\[
\begin{align*}
\dot{x} &= \sigma y - \alpha x \\
\dot{y} &= \rho x - y - xz + (y[n;t] - y)R \\
\dot{z} &= xy - \beta z
\end{align*}
\]
Here \( R \) is a coupling factor. This method works well from the standpoint of simplicity and ultimately from a circuit realization perspective. There have been circuit realizations proposed for the Lorenz equations\(^6\). In such circuits this coupling method is simply a resistive feed of a voltage across, or current into, a single arm.

### III. SEGMENT HOPPING

Now we turn our attention to the synthesis of the guide signal, \( y[n,t] \). Two new concepts are introduced here.

1. Continuous, digitally encoded chaos waveforms can be generated by piecing together the proper signal segments.
2. These segments can be bandlimited prior to storage, and can impress its characteristics upon a continuous time oscillator via synchronization.

The sequencing of stored segments in memory in some predefined order is commonly used in arbitrary waveform synthesis. What is new and not obvious about segment hopping is that an arbitrary controlled trajectory of a deterministic dynamical system can be generated, even though the system produces continuous and non-repeating trajectories in state space under the action of deterministic differential equations. This signal can furthermore be made to carry an arbitrary sequence of digital symbols representing encoded data. Thus the difference between this method and arbitrary waveform synthesis is that this method allows for the production of a signal carrying arbitrary data that appears to have been produced by the action of differential equations. This signal can furthermore be made to carry an arbitrary sequence of digital symbols representing encoded data. Thus the difference between this method and arbitrary waveform synthesis is that this method allows for the production of a signal carrying arbitrary data that appears to have been produced by the action of differential equations. This signal can furthermore be made to carry an arbitrary sequence of digital symbols representing encoded data. Thus the difference between this method and arbitrary waveform synthesis is that this method allows for the production of a signal carrying arbitrary data that appears to have been produced by the action of differential equations. This signal can furthermore be made to carry an arbitrary sequence of digital symbols representing encoded data. Thus the difference between this method and arbitrary waveform synthesis is that this method allows for the production of a signal carrying arbitrary data that appears to have been produced by the action of differential equations. This signal can furthermore be made to carry an arbitrary sequence of digital symbols representing encoded data. Thus the difference between this method and arbitrary waveform synthesis is that this method allows for the production of a signal carrying arbitrary data that appears to have been produced by the action of differential equations.

For example, in an 8-bit encoding there are 256 signal pieces, or segments, that can be put together to form any desired binary symbol sequence. The segments are assigned numbers from 0 to 255 according to the bit sequence they initiate.

In a physical implementation, a segment hopping system can be stored in a static memory device such as a ROM or EPROM and clocked out according to input the bit-sequence.

A more detailed description of the segment hopping process will be published shortly. In the following section we consider a complete computer model of the BCO synthesis method.

### IV. COMPUTER MODEL RESULTS

The first, important step is to determine the amount of bandlimiting that can be achieved. Any amount of filtering will cause some distortion to the waveforms. It is then important that the binary encoding remain preserved, and that the guide signal remain capable of synchronizing the continuous-time oscillator.

Given the determination of the average cycle time, \( T_{avp} \), as stated earlier, we applied a smooth, low-pass filter to the Lorenz oscillations having a cutoff-frequency at twice the average bit rate. With the filtering applied we stored the segments which source the appropriate 8-bit symbol sequences. Figs. 5 and 6 below show the filter characteristics and the response of the frequency spectra and the resultant attractor.

**Figure 5.** The result of the low-pass filtering of the state-variable of \( y(t) \).

**Figure 6.** Two-dimensional projection of a filtered Lorenz oscillation using the low-pass filter characteristic shown by its effect on the \( y \) spectrum in fig. 5.
The filtering is applied to all three state coordinates in order to produce a new attractor. It is with this new attractor that the symbolic dynamics of the system are determined and the associated segments are produced. Even though only one state-coordinate is used as a guide signal, namely y, the entire attractor must be transformed in order to synthesize the segments needed for segment hopping control.

In the following example we will encode and transmit the binary sequence, \( B = 11001101100010101111000110101100011 \). Fig. 7 shows what the desired output waveform \( x(t) \) would be for a typical Lorenz oscillator.

Given the bandlimited segments, we now can synthesize the guide signal, \( y[n] \). Fig. 8 is a plot of the frequency spectrum of the original, unfiltered transmit signal \( x(t) \), the result if \( x(t) \) was filtered using the low-pass filter described, and the frequency spectrum of the new transmit signal from the BCO synthesis method.

What is remarkable is how “flexible” the Lorenz oscillator can be. This is an important realization not only for this work but as we continue to work towards synthesizing signals even more compatible with traditional communication signal formats and standards. We have been able to demonstrate in the lab in prototypes that physical chaotic oscillators can be made to produce signals
that have constant timing intervals between Poincaré surface crossings by simply synchronizing it to an artificially synthesized waveform having those properties. This timing regularization is of utmost importance for use in commercially viable communication systems because it makes accurate clock timing recovery possible. Combined with bandwidth compression, as outlined here, we have solved some of the most critical technological problems for using chaotic oscillators in commercial systems.

V. CONCLUSION

What we have described here in a simple example is an enabling technology. In order to use chaotic systems and processes in commercial digital communication systems, the waveforms must be restricted in bandwidth in a controllable way. The focus must shift to the design of waveforms and sources, and away from the use of existing easily-constructed oscillators.

Our BCO synthesis technique is a simple, efficient, effective way to generate baseband signals for wireless communication. Some chaotic oscillators, such as the Colpitt’s circuit and the double scroll oscillator, produce signals with very different characteristics. However, this technique is a general procedure applicable to a variety of chaotic oscillators.

Another area of technology development that we are focusing on is in the development of oscillators that are more ideally suited for given communication channels. Although the Lorenz system has excellent properties for binary baseband digital signaling in white noise, there may be oscillations even more suitable that can produced with efficient circuitry.

Since Ott, Grebogi, and Yorke’s initial formalism for controlling chaotic processes using small perturbations, there has been a consistent march toward development of technological applications. Arguably, the application to communications technology has held the most promise and has been the source of the most activity. Beyond simple applications is the realm of commercially viable applications that have characteristics that make chaotic dynamics technology attractive to the fast-moving world of the telecommunications industry.

Some of the improvement areas that make chaotic dynamics technology attractive are increases in efficiency, reduction of system complexity, increase in transmission ranges and digital transmission data rates.

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VII. REFERENCES